

尚途

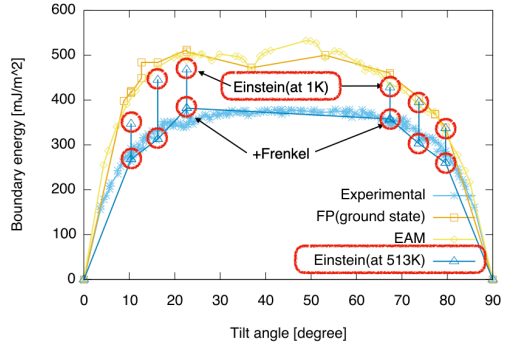
$dE = Ed - nE_0$ ① 他? 女陷

$dF(T) = F_d(T) - nF(T)$ ② 非調和

$F = E_0 + F_{\text{Einstein}} + F_{\text{Frenkel}}$ ③ FP

調和 非 ~ 0 (513K)

“Finite-temperature first-principles calculations of Al<100> symmetric tilt grain-boundary energy”



S. R. Nishitani, Phil. Mag., vol. 101, pp. 622–42, 2021, (https://doi.org/10.1080/14786435.2020.1855371).

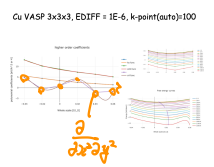
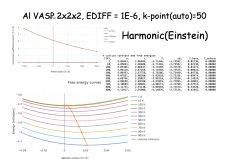
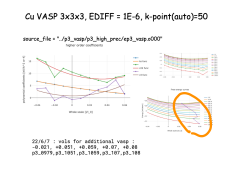
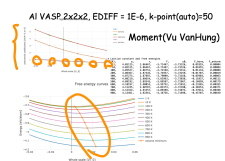
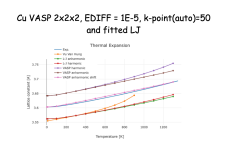
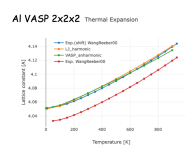
Finite Temp VASP Al<100> tilt GB

Static/Analytical Moment method

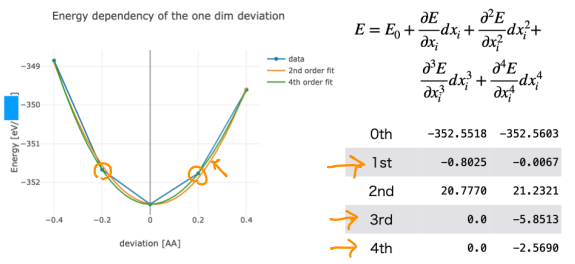
- thermal expansion

Al

Cu



SMM (Statistical Moment Method)



N. Tan, V. V. Hung, “Investigation of the Thermodynamic Properties of Anharmonic Crystals by the Momentum Method”, Phys. Stat. Sol. (B), vol. 149, pp.511–519, 1988.

$F_{\text{anharmonic}} \approx E_0 + F_{\text{Einstein}}$

$\theta = k_B T, x = \frac{\hbar\omega}{2\theta}$

$F_0 = \sum \sigma(\omega_j)$

$F_{\text{Einstein}} = \theta \left[x + \log(1 - e^{-2x}) \right]$

$k = \frac{1}{2} \sum \frac{\partial^2 \psi_{ij}}{\partial x_i^2} = m\omega^2$

$\gamma_i = \frac{1}{48} \sum \frac{\partial^3 \psi_{ij}}{\partial x_i^3}, \gamma_z = \frac{6}{48} \sum \frac{\partial^3 \psi_{ij}}{\partial x_i^3}$

$F_{\text{higher}}(T) = E_d(T) + \sum_{j=1,2,3} h_{ij}(T) \times h_j^* \times h_j^* \times h_j^*$

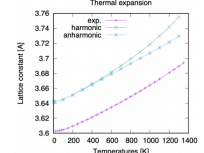
PAW - PBE

TABLE 1. Higher order coefficients of the energy potential.

indices	analytical	numerical	numerical harmonic	VASP(2x2x2)	VASP(3x3x3)
0 [0 0 0]	-2.47640	-2.47640	-2.47640	-2.47640	-2.47640
1 [1 0 0]	6.60322	6.60486	6.60288	6.60320	6.60320
2 [2 0 0]	6.60322	6.60486	6.60288	6.60320	6.60320
3 [3 0 0]	6.60322	6.60486	6.60288	6.60320	6.60320
4 [4 0 0]	1.7206-12	2.8406-13	0.00000	0.00000	0.00000
5 [1 1 0]	1.6006-11	1.126-12	0.00000	0.00000	0.00000
6 [1 1 1]	1.7206-12	2.8406-13	0.00000	0.00000	0.00000
7 [2 1 0]	-1.126-10	-1.126-10	0.00000	0.00000	0.00000
8 [2 1 1]	-1.126-10	-1.126-10	0.00000	0.00000	0.00000
9 [3 1 0]	-2.1476-09	2.406-11	0.00000	0.00000	0.00000
10 [3 1 1]	-2.1476-09	2.406-11	0.00000	0.00000	0.00000
11 [2 2 0]	29.14019	29.27755	29.34444	29.34444	29.34444
12 [2 2 1]	29.14019	29.27755	29.34444	29.34444	29.34444
13 [3 2 0]	6.25142	6.51089	6.84033	6.84033	6.84033
14 [3 2 1]	6.25142	6.51089	6.84033	6.84033	6.84033
15 [4 2 0]	6.25142	6.51089	6.84033	6.84033	6.84033
standard dev.	0.000000	0.000014	0.000025	0.000007	0.000001

- Analytical/numerical anharmonic free energy of pure Cu or Al system
- Conclusions
 - VuVanHung's Moment method
 - Thermal expansion
 - Cu (x), Al (o)
 - Frenkel Integration
 - $F_{\text{Frenkel}} \neq F_{\text{moment}}$
 - $F_{\text{VASP}} = F_{\text{moment}} + \int \frac{dE}{d\lambda} d\lambda \neq F_{\text{Einstein}} + F_{\text{Frenkel Integral}}$

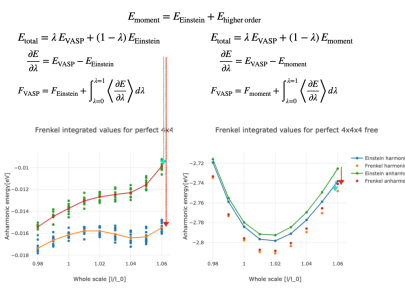
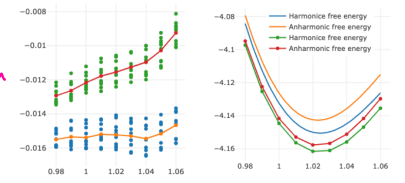
Cu 完全結晶の解析的・数值的 非調和 自由エネルギー



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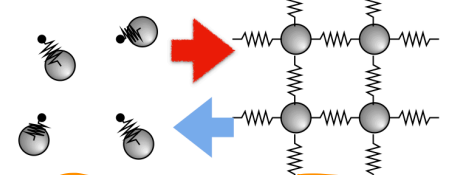
outline

- Finite Temperature First Principles Calculations
 - Al <100> symmetric tilt boundary energy
- How to include anharmonicity
 - analytical : SMM (Statistical Moment Method)
 - numerical : Frenkel (Monte Carlo based Thermodynamic Integration)
- Al and Cu system
- VASP vs fitted LJ (2nd NN)



Dynamic/Numerical Frenkel(MC-TI)

Einstein model VASP model



$F_i = E_i^0 - k_B T \ln Z_i$

$E_i^0 = E_i^0 - k_B T \sum_{j=x,y,z} \ln \left(\frac{\exp(-\hbar\omega_j/2k_B T)}{1 - \exp(-\hbar\omega_j/k_B T)} \right)$

$E_{\text{total}} = \lambda E_{\text{VASP}} + (1 - \lambda) E_{\text{Einstein}}$

$\frac{\partial E}{\partial \lambda} = E_{\text{VASP}} - E_{\text{Einstein}}$

$F_{\text{VASP}} = F_{\text{Einstein}} + \int_{\lambda=0}^{\lambda=1} \left\langle \frac{\partial E}{\partial \lambda} \right\rangle d\lambda$

Al

Cu

skipped